Validation of Special Sensor Ultraviolet Limb Imager (SSULI) Ionospheric Tomography using ALTAIR Incoherent Scatter Radar Measurements

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Introduction

Data assimilation models are used to specify and forecast ionospheric conditions

- These models rely heavily on data accuracy
- Requires validation to ensure consistency between measurement sets
- Requires assessment of measurement strengths and weaknesses
- UV limb-scanner measurements of emissions that originate in the ionosphere are useful for specifying the latitude-altitude distribution of plasma along the orbit plane
 - Difficult to validate because measurements for direct comparison do not exist
- What are we trying to learn?
 - How well do the UV measurements capture the underlying distribution of plasma in the ionosphere?
 - How are the measurements affected by illumination and ionospheric gradients?
 - How well do the UV measurements compare to other measurement techniques?
- We used coincident measurements of the latitude-altitude distribution of electrons using the incoherent scatter radar at ALTAIR to assess how accurately UV measurements can specify the ionospheric plasma distribution.





SSULI Measurement Scenario







Overview

- What are we trying to do?
 - Specific application: Demonstrate and validate on-orbit specification of the ionosphere
 - Approach: Use aggregates of limb scan information to infer the 2-D distribution of O⁺ ions in the ionosphere
- Brightness measurements are linear in the volume emission rate
 - Analogous to Computerized lonospheric Tomography → linear in the electron density
 - Noise on brightness measurements obeys Poisson statistics – not the Normal Distribution



Volume emission rate, ε : $\varepsilon(z,\lambda,\phi) = \alpha n_{e}(z,\lambda,\phi) n_{o^{+}}(z,\lambda,\phi)$

$$4\pi I = 10^{-6} \int_0^\infty \varepsilon(s, z, \lambda, \phi) \exp(-\tau(s, z, \lambda, \phi)) ds(z, \lambda, \phi)$$





O I 911 Å Emission & Absorption

- The 911 Å emission is only excited by radiative recombination of Fregion O⁺ ions and electrons:
 - $O^+ + e^- \rightarrow O + hv (911 \text{ Å})$
 - Rate coefficient: α = 3.5×10⁻¹³ (1160./T[K])^{1/2} cm⁻³ s⁻¹ (Melendez-Alvira et al, 1999)
- The 911 Å emission is attenuated by atomic oxygen, molecular oxygen, and molecular nitrogen:
 - O: Photoionization (Conway, 1989, scaled)
 - O + hv (911 Å) \rightarrow O⁺ + e⁻
 - Cross-section: $\sigma = 3.93 \times 10^{-18} \text{ cm}^2$
 - This has not been previously identified as a loss process
 - O₂: Photoionization & absorption (Conway, 1989)
 - O₂ + hν (~900 Å) → O⁺ + O + e⁻
 - $O_2^- + hv (\sim 900 \text{ Å}) \rightarrow O_2^*$
 - Cross-section: $\sigma = 15.34 \times 10^{-18} \text{ cm}^2$
 - N₂: Absorption (Kirby et al., 1979)
 - $\text{ N}_2 + h\nu \ (\text{~900 Å}) \rightarrow \text{N}_2^*$
 - Cross-section: $\sigma = 14.50 \times 10^{-18} \text{ cm}^2$







Total Cross-sections for O, O₂ & N₂ near 911 Å



For Comparison: Peak of Schumann-Runge O₂ Absorption ~14.8 MBarn





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Factors to Consider

- Inversion of 911 Å measurements requires:
 - Accurate calibration
 - Accurate model of the measurements
 - Line-of-sight integration (quadrature scheme)
 - Accurate physics: absorption, photochemical process
 - Statistical representation of measurement noise
- Calibration
 - Cannot use stars to calibrate at 911 Å because there are no stellar photons that reach the Earth due to interstellar absorption
 - Affects the magnitude of the retrieved density, not the morphology
- Model of the measurements
 - Measurement model affects the morphology, not the magnitude
 - Quadrature scheme needs to reproduce expected variations
 - Measurement statistics guide where the algorithm attributes
 - emission
 - Absorption determines where the emission can be attributed
 - Photochemical process converts the volume emission rate to the product: electron density







Interstellar Absorption

- There is essentially zero stellar flux at wavelengths shorter than 912 Å
 - There is flux at much shorter wavelengths
 - Also, dwarf stars in the solar neighborhood can be seen, but their fluxes are low
- Figure at left shows a stellar spectrum used to calibrate the SSULI instruments at longer wavelengths
 - Taken from: Morales et al., "Farultraviolet absolute flux of α Virginis", The Astrophysical Journal, 530:403-407, (2000)
- ➢ How do we calibrate SSULI at 911 Å?
 - We use ground truth radar measurements.





Interstellar absorption Cut-off ~920 Å





The Approach

Because UV measurements for comparison do not exist

- Need to validate measurements by deriving a product that can be measured by an alternative means
- This approach provides an end-to-end test from the measurements through the interpretation process
 - Assessment of calibration
 - Assessment of observation scenario
- ALTAIR Incoherent Scatter Radar
 - Measures two-dimensional electron density distribution along the orbit plane
 - Beneath the Equatorial Ionization Anomaly
 - High densities
 - Structure and gradients







Ionospheric Tomography Algorithm

$$4\pi I = 10^{-6} \sum_{i} \varepsilon(z,\lambda,\phi) \Big[\exp(-\tau(s,z,\lambda,\phi)) \Delta s_{i}(z,\lambda,\phi) \Big]$$

- Line-of-sight integrals are replaced by summations assuming constant volume emission rate in a voxel weighted by optical extinction
- The result is a large sparse linear system of equations
- This system is solved using the iterative Richardson-Lucy algorithm
 - Non-negative
 - Tailored to Poisson random deviates
- Solution physicality is ensured by regularizing to a partial differential equation: The diffusion equation



Ax = b

$$x_{j+1} = x_j \otimes \frac{A^{T}}{A^{T}(\bar{1})} \left(\frac{b}{Ax_j}\right)$$

$$\frac{\partial n}{\partial t} = \nabla \Box \left(\overline{D} \nabla n \right) \Longrightarrow 0 = \nabla^2 n$$



Scatter Plots: Optically Thick and Thin August 19, 2014

- Plotted scatter plots of corrected SSULI data versus ALTAIR data
 - Dashed line is the unity slope line indicating perfect agreement
- Calibration scale factor determined for each inversion
- Top: scatter plot without absorption
 - SSULI overestimating highest densities by ~50% when density is near 7×10⁵ cm⁻³
- Bottom: scatter plot with re-ionization of O
 - Better agreement at all densities
 - SSULI overestimating highest densities by ~30% when density is near 7×10⁵ cm⁻³
 - Scatter of distribution is tighter than it is without the re-ionization





Scatter Plots: Pure Absorption August 19, 2014



Model with absorption by O & O_2 has lowest scatter





Scatter Plots: Optically Thick and Thin October 12, 2014

- Plotted scatter plots of corrected SSULI data versus ALTAIR data
 - Dashed line is the unity slope line indicating perfect agreement
- Calibration scale factor determined for each inversion
- Top: scatter plot without absorption
 - Correlation is good, but there is an outlier population
- Bottom: scatter plot with re-ionization of O & O₂ absorption
 - Better agreement at all densities
 - Scatter of distribution is tighter than it is without the re-ionization
 - Outlier population is significantly reduced









SSULI/ALTAIR: August 19, 2014 -Dusk Pass-







SSULI/ALTAIR: September 4, 2014 -Dusk Pass-



No absorption





SSULI/ALTAIR: October 12, 2014 -Dusk Pass-



No absorption





SSULI/ALTAIR: September 29, 2014 -Dawn Pass-



No absorption

5/22/2015



With absorption



altitudes

Summary

- We compared the results of UV tomography using UV measurements made by the SSULI sensor to ALTAIR
 - Excellent agreement with the altitude/latitude distributions from the two measurements for the dusk passes
 - Dawn passes are still under investigation
 - The measurements were made in the terminator region, which are typically not used because they are difficult to interpret
- Our analysis approach entailed
 - New iterative Image Space Reconstruction Algorithm -- Richardson-Lucy technique -- handles Poisson noise explicitly and is non-negative
 Can work on data with very low signal-to-noise ratio
 - Physicality constraint using regularization to the isotropic diffusion equation
 - Inclusion of re-ionization of O & absorption by O₂ by the 911 Å emission – found to be important by this analysis







Conclusions

- Tomography approach produces accurate electron density distribution
 - Calibration of the UV data determined during the process
 - Good agreement between tomography products and ALTAIR measurements validates the measurements → good for use by GAIM model
- Neutral absorption is an important consideration for interpretation of the UV measurements
 - Results are improved when absorption by O and O_2 are included
 - When absorption by N₂ is included in the model, the agreement with ALTAIR is degraded
 - Suggests that N_2 cross-sections near 900 Å need further investigation







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Physicality Constraint

- Regularization to a differential equation is an approach used in the computer graphics modeling community
 - Improves computer rendering by generating a smooth surface from facet information
- We use the time independent diffusion equation

$$\frac{\partial n}{\partial t} = \nabla \Box \left(\overline{D} \nabla n \right) \Longrightarrow 0 = \nabla^2 n \quad (time \ independent)$$

- Currently, we assume uniform, isotropic transport
 - Permits the algorithms to produce reasonable results during daytime and at night
 - Will work for either ionospheric emissions (nighttime ionosphere) or for emission generated by neutral species (O and N₂ in the dayglow)
 - However, some emissions, for example O I 1356 Å, have both ionospheric and thermospheric components during the daytime
 - Drives eventual need for non-isotropic, non-uniform diffusion approximation
- Implemented using the Successive Over-Relaxation approximation
 - Makes small steps to "relax" solution to the diffusion approximation





Successive Over-Relaxation (SOR)

- > We chose this iterative approach to solve the diffusion equation
 - Desired a method with low computational overhead
 - Wanted a means to guide the algorithms to a physically meaningful solution
- > Approximating the diffusion equation at time step k+1 by finite difference equations (assuming $\Delta x = \Delta y$, i & j are cell indices):

$$n_{i,j}^{k+1} = n_{i,j}^{k} - \frac{D\Delta t}{\left(\Delta x\right)^{2}} \left(n_{i-1,j}^{k} + n_{i+1,j}^{k} + n_{i,j-1}^{k} + n_{i,j+1}^{k} - 4n_{i,j}^{k}\right)$$

To ensure a stable solution, the maximum time step size allowed is limited by the diffusion time across the cell:

$$W \equiv \frac{D\Delta t}{\left(\Delta x\right)^2} \le \frac{1}{4}$$

• We refer to W as the diffusion weight and use it to tune the weighting of the physicality constraint



